

T production in dAu collisions at RHIC

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In collaboration with E. G. Ferreiro, F. Fleuret, J. P. Lansberg and A. Rakotozafindrabe

Outline

- Introduction and motivations
- Experimental situation
- \bigcirc On the kinematics of Υ production
- The Glauber Monte Carlo
- 5 Results for dAu collisions for Υ
- 6 EMC effect for gluons
- Conclusions and perspectives

 Extend to ↑ the study of CNM effects (shadowing + absorption) on production of quarkonia

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- Two main production schemes $(2 \rightarrow 1, 2 \rightarrow 2)$

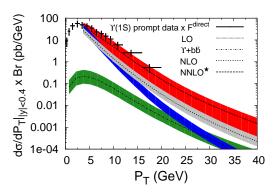
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- Comparison of three differents shadowing parametrisations
- Three absorption cross sections

P. Artoisenet, J. Campbell, J.P. Lansberg, F. Maltoni, Phys. Rev. Lett. 101, 152001 (2008).

D. Acosta et al. (CDF collaboration), Phys. Rev. Lett 88, 161802 (2002).

Results at 1.8 TeV



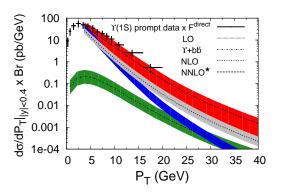
CSM describes well the data at NNLO*



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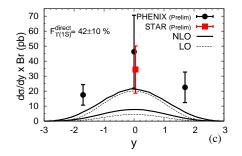


- CSM describes well the data at NNLO*
- However LO CSM is sufficient to describe low pT data

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S. J. Brodsky and J. P. Lansberg, Phys. Rev. D81, 014004 (2010).
P. Djawotho et al., J. Phys. G34, s947 (2007); T. Ullrich (private communication) (STAR)
C.L. da Silva, Nucl. Phys. A830, 227c (2009); L.L. Levy, Nucl. Phys. A830, 353c (2009); W. Xie et al., J. Phys. A774, 693 (2006) (PHENIX)

Results at 200 GeV



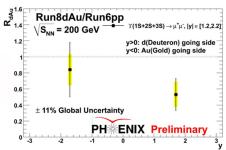
- Upper dashed line, $m_b = 4.5$ GeV, $\mu_R = M_T$, $\mu_F = 2M_T$
- Lower dashed line, $m_b = 5.0$ GeV, $\mu_R = 2M_T$, $\mu_F = M_T$

We take the parameters of the upper curve in the following.

H. Liua and the STAR collaboration, Nucl. Phys. A830, 235c (2009)
H. Pereira Da Costa for the PHENIX collaboration, talk at the rencontres de Moriond, March 15, 2010

3 data points :

- ullet -2.2 < y < -1.2: PHENIX $R_{dAu} = 0.84 \pm 0.34$ (stat.) ± 0.28 (sys.)
- ullet 1.2 < y < 2.2: PHENIX $R_{dAu} = 0.53 \pm 0.20$ (stat.) \pm 0.16 (sys.)



ullet |y| < 0.5: STAR - $R_{dAu} = 0.98 \pm 0.32$ (stat.) \pm 0.28 (sys.)

If $\mathcal{F}_g^A(x, \vec{r}, z, \mu_f)$ gives the distribution of a gluon of mom. fract. x at a position \vec{r}, z in a nucleus A, the differential cross-section reads:

$$\frac{d\sigma_{AB}}{dy\ dP_T\ d\vec{b}} =$$

$$2
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 kinematics with intrinsic $ho_{\mathcal{T}}$

$$2 \rightarrow 2$$
 kinematics with extrinsic p_T

If $\mathcal{F}_{g}^{A}(x,\vec{r},z,\mu_{f})$ gives the distribution of a gluon of mom. fract. x at a position \vec{r} , z in a nucleus A, the differential cross-section reads:

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 $2 \rightarrow 1$ kinematics with intrinsic p_T

$$\int d\vec{r}_{A} dz_{A} dz_{B}$$

$$\times \mathcal{F}_{g}^{A}(\mathbf{x}_{1}^{0}, \vec{r}_{A}, z_{A}, \mu_{f}) \mathcal{F}_{g}^{B}(\mathbf{x}_{2}^{0}, \vec{r}_{B}, z_{B}, \mu_{f})$$

$$\times \mathcal{F}_{g}^{A}(\mathbf{x}_{1}, \vec{r}_{A}, z_{A}, \mu_{f}) \mathcal{F}_{g}^{B}(\mathbf{x}_{2}, \vec{r}_{B}, z_{B}, \mu_{f})$$

$$\times \sigma_{gg}^{\text{Intr.}}(\mathbf{x}_{1}^{0}, \mathbf{x}_{2}^{0})$$

$$\times S_{A}(\vec{r}_{A}, z_{A}) S_{B}(\vec{r}_{B}, z_{B})$$

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 $2 \rightarrow 2$ kinematics with extrinsic p_T

$$\int d\mathbf{x}_{1}d\mathbf{x}_{2} \int d\vec{r}_{A}dz_{A}dz_{B}
\times \mathcal{F}_{g}^{A}(\mathbf{x}_{1}, \vec{r}_{A}, z_{A}, \mu_{f}) \mathcal{F}_{g}^{B}(\mathbf{x}_{2}, \vec{r}_{B}, z_{B}, \mu_{f})
\times 2\hat{\mathbf{s}} P_{T} \frac{d\sigma_{gg \to \Upsilon + g}}{d\hat{\tau}} \delta(\hat{\mathbf{s}} - \hat{\mathbf{t}} - \hat{\mathbf{u}} - M^{2})
\times S_{A}(\vec{r}, z_{A}) S_{B}(\vec{r}_{B}, z_{B})$$

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$$\times \mathcal{F}_{g}^{A}(\mathbf{x}_{1}^{1}, \vec{r}_{A}, z_{A}, \mu_{f}) \mathcal{F}_{g}^{B}(\mathbf{x}_{2}^{0}, \vec{r}_{B}, z_{B}, \mu_{f})$$

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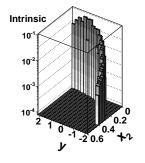
$$\times \mathcal{F}_{g}^{A}(\mathbf{x}_{1}, \mathbf{x}_{A}, \mu_{f}) \mathcal{F}_{g}^{B}(\mathbf{x}_{2}, \mathbf{x}_{A}, \mu_{f}) \mathcal{F}_{g}^{B}(\mathbf{x}_{2}, \mathbf{x}_{A}, \mu_{f}) \mathcal{F}_{g}^{B}(\mathbf{x}_{2}, \mu_{f}) \mathcal{F}_{g}^{B}(\mathbf{x}$$

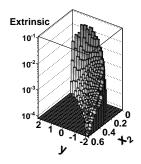
$$x_{1,2} = \frac{m_T}{\sqrt{s_{NN}}} \exp(\pm y) \equiv x_{1,2}^0(y, P_T)$$

 $2 \rightarrow 2$ kinematics with extrinsic p_T

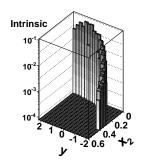
$$\begin{split} &\int\!\! d\mathbf{x}_1 d\mathbf{x}_2 \int d\vec{r}_A dz_A dz_B \\ &\times \mathcal{F}_g^A(\mathbf{x}_1, \vec{r}_A, z_A, \mu_f) \mathcal{F}_g^B(\mathbf{x}_2, \vec{r}_B, z_B, \mu_f) \\ &\times 2 \mathbf{\hat{s}} P_T \frac{d\sigma_{gg \to \Upsilon + g}}{d\mathbf{\hat{t}}} \delta(\mathbf{\hat{s}} - \mathbf{\hat{t}} - \mathbf{\hat{u}} - \mathbf{M}^2) \\ &\times S_A(\vec{r}, z_A) S_B(\vec{r}_B, z_B) \end{split}$$

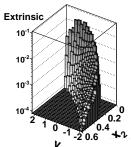
$$\delta(..) \rightarrow x_2 = \frac{x_1 m_T \sqrt{s_{NN}} e^{-y} - M^2}{\sqrt{s_{NN}} (\sqrt{s_{NN}} x_1 - m_T e^y)}$$





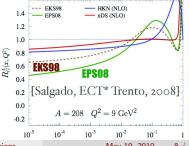
For a given couple (y, p_T) , x_2 is larger in the extrinsic scheme

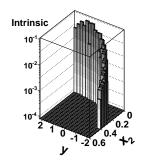


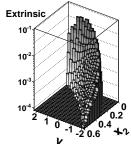


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Antishadowing peak at $\sim 10^{-1}$



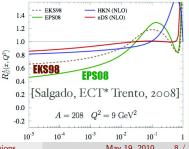




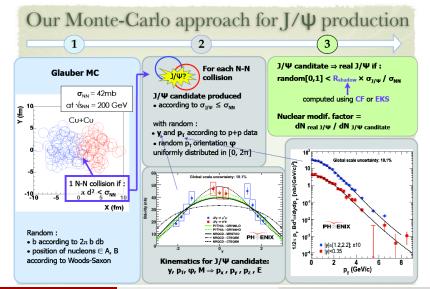
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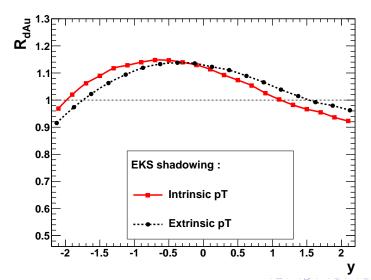
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We expect different shadowing effects in both cases.

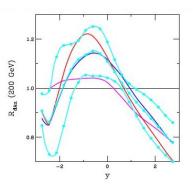


The Glauber Monte Carlo (for Υ here)





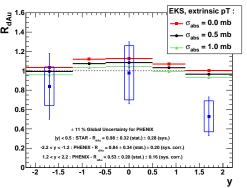
R. Vogt, talk at Joint CATHIE-TECHQM Meeting, BNL, December 14-18, 2009



EKS98 (blue), nDSG (magenta), EPS08 (red), EPS09 (cyan)

- One has to be careful about binning effect (usually decrease the modifications)
- Interesting to see the difference between 2->1 (as done by R. Vogt) and 2->2 (kinematics for LO CSM)

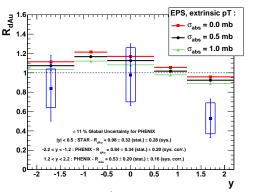
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- backward: ok within uncertainties;
- central: reasonable job $R_{dAu} > 1$ (for any σ_{abs});
- forward: clearly too high (for any σ_{abs});



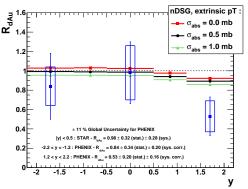
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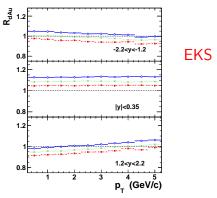
- backward: slightly too high (but ok within uncertainties);
- central: reasonable job $R_{dAu} > 1$ (for any σ_{abs});
- forward: clearly too high (for any σ_{abs}), though 'better' than EKS;

Results for dAu collisions for Υ

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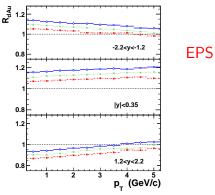
- backward: ok within uncertainties;
- central: reasonable job R_{dAu} 1;
- forward : clearly too high (for any σ_{abs}), though 'better' than EKS and EPS:



The extrinsic scheme enables to predict the pT dependence, which is non trivial

- In blue, $\sigma_{abs} = 0.0 \text{ mb}$
- In green, $\sigma_{abs} = 0.5 \text{ mb}$
- In red, $\sigma_{abs} = 1.0 \text{ mb}$

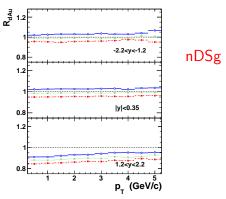




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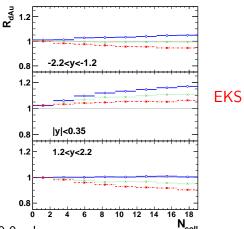




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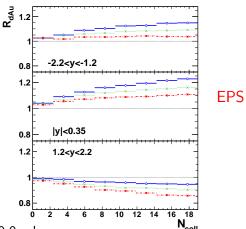
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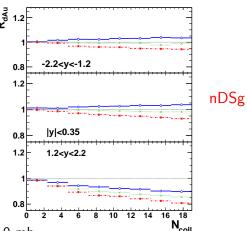
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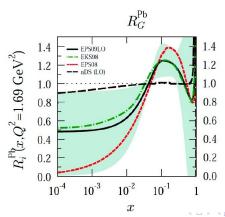


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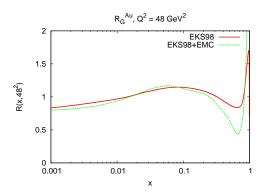
EMC effect for gluons

- Tension between the theory and the PHENIX point in the backward region
- The backward region correspond to the EMC region (x > 0.1)
- EMC effect basically unknown for the gluon



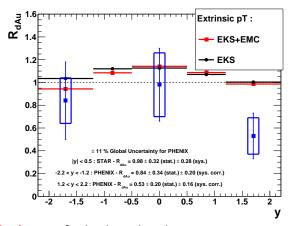
EMC effect for gluons

- Let us try to increase the suppression of g(x) in the EMC region
- Keeping momentum conservation : $\int xg(x) dx = Cst$



EMC effect for gluons

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Works better for backward region

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- Within the commonly accepted σ_{abs} , one should expect an excess of Υ
- ... unless there is no antishadowing (see nDSg)

